# Design Optimization of Soft-Switched Insulated DC/DC Converters With Active Voltage Clamp 

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#### Abstract

The use of an active voltage clamp in insulated DC/DC converters allows a limitation of the voltage stress due to transformer leakage inductance, a reduction of the commutation losses, and an increase of the switching frequency. It is therefore important to ensure soft switching in the whole operating range, penalty a considerable worsening of converter efficiency and electromagnetic noise generation.

In this paper a design approach of insulated converters with active voltage clamp is presented, which allows soft switching from no load to full load. The approach is described for the case of a flyback converter but is valid also for Sepic and Cuk topologies.

Experimental results of a $100 \mathrm{~W}-300 \mathrm{kHz}$ converter demonstrate the excellent performances obtained according to the proposed design procedure: switch overvoltage less than $\mathbf{2 0 \%}$ of the nominal value in any operating condition and overall converter efficiency above $\mathbf{8 3 \%}$ from $50 \%$ to $100 \%$ of rated load.


## I. InTRODUCTION

When electric insulation is a prime concern in dc-dc conversion, flyback, Sepic and Cuk topologies are good candidates for their relative simplicity as compared to other solutions. The main problem of these converters is represented by the transformer leakage inductance, which may cause severe voltage stress on the main devices and contributes significantly to conducted and radiated EMI.

Passive clamps, like the usual R-C-D snubber circuit, help to limit the switch overvoltage at the expense of higher losses, while damping of the high frequency parasitic oscillations calls for additional R-C snubbers, further decreasing the overall efficiency.

Active clamps, consisting of an auxiliary switch in series to a clamp capacitor, provide a viable solution to reduce the voltage stress without increasing appreciably the converter losses. Moreover, the auxiliary switch of the active clamp allows a soft commutation of the main devices [1-3]. For this purpose, the energy stored in the transformer is exploited to discharge the switch parasitic capacitance before turning on, thus achieving zero-voltage
switching of the main switch, soft turn off of the freewheeling diode and elimination of high frequency parasitic oscillations.

Previous works on this topic have demonstrated that soft switching is feasible for both discontinuous [1] and continuous [23] conduction mode.

The purpose of this work is twofold: first, to derive design criteria of a flyback converter with active clamp soft switched in the whole load range; second, to show that the results obtained for the flyback converter can be directly extended also to Sepic and Cuk topologies.

The theoretical forecasts are experimentally tested on a 100 W flyback converter switching at 300 kHz , which in fact demonstrates high efficiency, small voltage stresses and no-load to full-load regulation under soft-switched conditions.


Fig. 1 - Flyback converter with active clamp

## II. Analysis of Flyback Converter with Active Clamp - a Review

In this section, the behavior of flyback converter with active clamp is reviewed both for Continuous (CCM) and Discontinuous (DCM) Conduction Modes, while in the next section it will be shown that the same analysis is also valid for Sepic and Cuk converters.

Fig. 1 shows the basic scheme of a flyback converter with active clamp. Inductance $\mathrm{L}_{\mathrm{R}}$ includes the transformer leakage inductance
while $C_{P}$ accounts for the switch output capacitances. In the following analysis, ideal switching components are assumed.


Fig. 2 - Main converter waveforms in CCM operation

a)

c)

e)

b)

d)

f)

Fig. 3 - Subtopologies corresponding to different switch states

## .A. Continuous Conduction Mode (CCM)

As it can be seen in Fig. 2, which shows the main converter waveforms in CCM, each switching period is subdivided into six intervals, each of them correspondent to a different topology (see Fig.3). In order to simplify the analysis, the relative time of each interval is set to zero at the beginning of the interval.

1) Interval $t_{0-t}{ }_{1}$ (Fig.3a): At instant $\mathrm{t}_{0}$ the freewheeling diode D stops conduction and switch $S_{1}$ is on while switch $S_{2}$ is off. The currents in leakage inductance $\mathrm{L}_{\mathrm{R}}$ and magnetizing inductance $\mathrm{L}_{\mathrm{M}}$ are equal and increase linearly with a slope dependent on the input voltage:

$$
\begin{equation*}
\mathrm{i}_{\mathrm{M}}(\mathrm{t})=\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{0}\right)+\frac{\mathrm{V}_{\mathrm{i}}}{\mathrm{~L}_{\mathrm{M}}(1+\beta)} \cdot \mathrm{t} \tag{1}
\end{equation*}
$$

where $\beta=\frac{L_{R}}{L_{M}}$.
2) Interval $t_{1}^{-t_{2}}$ (Fig.3b): $S_{1}$ is turned off at instant $t_{1}$ and $C_{P}$ is then charged almost linearly by the magnetizing inductance current:

$$
\begin{equation*}
\mathrm{v}_{\mathrm{C}_{\mathrm{p}}}(\mathrm{t})=\frac{\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right)}{\mathrm{C}_{\mathrm{P}}} \cdot \mathrm{t} \tag{2}
\end{equation*}
$$

At instant $t_{2}$, this voltage equals $\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{C}}$ and diode $\mathrm{D}_{2}$ starts conducting.
3) Interval $t_{2}-t_{3}$ (Fig.3c): After conduction of $\mathrm{D}_{2}$, inductors $\mathrm{L}_{\mathrm{R}}$ and $\mathrm{L}_{\mathrm{M}}$ resonate with capacitors $\mathrm{C}_{\mathrm{R}}$ and $\mathrm{C}_{\mathrm{P}}$ until the transformer secondary voltage becomes greater than the output voltage and the freewheeling diode starts conducting (instant $\mathrm{t}_{3}$ ). Resonant current and voltage are given by the following equations:

$$
\left\{\begin{array}{l}
i_{M}(t)=i_{R}(t)=A_{1} \cdot \sin \left(\omega_{1} t+\varphi_{1}\right)  \tag{3.a}\\
v_{C_{R}}(t)=-A_{1} Z_{1} \cdot \cos \left(\omega_{1} t+\varphi_{1}\right)
\end{array}\right.
$$

where

$$
\begin{align*}
& \left\{\begin{array}{l}
A_{1}=\sqrt{\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right)^{2}+\left(\frac{\mathrm{v}_{\mathrm{C}_{\mathrm{R} 0}}}{\mathrm{Z}_{1}}\right)^{2}} \\
\operatorname{tg}\left(\varphi_{1}\right)=\frac{\mathrm{Z}_{1} \cdot \mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right)}{-\mathrm{v}_{\mathrm{C}_{\mathrm{R} 0}}}
\end{array}\right.  \tag{3.b}\\
& \omega_{1}=\frac{1}{\sqrt{\left(\mathrm{C}_{\mathrm{R}}+\mathrm{C}_{\mathrm{P}}\right) \cdot \mathrm{L}_{\mathrm{M}}(1+\beta)}}  \tag{3.c}\\
& Z_{1}=\sqrt{\frac{\mathrm{L}_{\mathrm{M}}(1+\beta)}{\mathrm{C}_{\mathrm{R}}+\mathrm{C}_{\mathrm{P}}}} \tag{3.d}
\end{align*}
$$

and ${ }^{\mathrm{V}_{\mathrm{R} 0}}$ is the initial value of the voltage across $\mathrm{C}_{\mathrm{R}}$. This interval ends when

$$
\begin{equation*}
\mathrm{v}_{\mathrm{C}_{\mathrm{R}}}\left(\mathrm{t}_{3}\right)=\mathrm{V}_{\mathrm{op}} \cdot(1+\beta) \tag{4}
\end{equation*}
$$

where $\mathrm{V}_{\mathrm{op}}=\mathrm{V}_{\mathrm{o}} / \mathrm{n}$ is the output voltage reported to the primary side.
4) Interval $t_{3^{-t}}^{4}$ (Fig.3d): During this interval the magnetizing current decreases linearly, transferring energy to the output, while
$L_{R}$ resonates with $C_{R}$. Thus from the analysis of Fig.3d we can write:

$$
\begin{align*}
& \mathrm{i}_{\mathrm{M}}(\mathrm{t})=\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{3}\right)-\frac{\mathrm{V}_{\mathrm{op}}}{\mathrm{~L}_{\mathrm{M}}} \cdot \mathrm{t}  \tag{5}\\
& \left\{\begin{array}{l}
\mathrm{i}_{\mathrm{R}}(\mathrm{t})=\mathrm{A}_{2} \cdot \sin \left(\omega_{2} \mathrm{t}+\varphi_{2}\right) \\
\mathrm{v}_{\mathrm{C}_{\mathrm{R}}}(\mathrm{t})=\mathrm{V}_{\mathrm{op}}-\mathrm{A}_{2} \mathrm{Z}_{2} \cdot \cos \left(\omega_{2} \mathrm{t}+\varphi_{2}\right) \\
\mathrm{A}_{2}=\sqrt{\mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{3}\right)^{2}+\left(\frac{\mathrm{V}_{\mathrm{op}} \beta}{\mathrm{Z}_{2}}\right)^{2}}, \mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{3}\right)=\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{3}\right) \\
\operatorname{tg}\left(\varphi_{2}\right)=\frac{\mathrm{Z}_{2} \cdot \mathrm{i}_{R}\left(\mathrm{t}_{3}\right)}{-V_{o p} \beta}
\end{array}\right.  \tag{6.a}\\
& \omega_{2}=\frac{1}{\sqrt{\left(\mathrm{C}_{R}+C_{P}\right) \cdot L_{R}}}  \tag{6.b}\\
& \mathrm{Z}_{2}=\sqrt{\frac{L_{R}}{\mathrm{C}_{R}+C_{P}}} \tag{6.c}
\end{align*}
$$

Note that the auxiliary switch $S_{2}$ is turned on in lossless manner before the clamp capacitor current reverses.
5) Interval $t_{4}{ }^{-t} 5$ (Fig.3e): At instant $t_{4} S_{2}$ is turned off and $L_{R}$ resonates with $\mathrm{C}_{\mathrm{P}}$ bringing its voltage to zero (assuming that it has enough energy) and forcing the conduction of $\mathrm{D}_{1}$ (instant $\mathrm{t}_{5}$ ). Current and voltage behavior is described by the equations:

$$
\begin{align*}
& \left\{\begin{array}{l}
\mathrm{i}_{\mathrm{R}}(\mathrm{t})=\mathrm{A}_{3} \cdot \sin \left(\omega_{3} \mathrm{t}+\varphi_{3}\right) \\
\mathrm{v}_{\mathrm{C}_{\mathrm{p}}}(\mathrm{t})=\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}-\mathrm{A}_{3} \mathrm{Z}_{3} \cdot \cos \left(\omega_{3} \mathrm{t}+\varphi_{3}\right)
\end{array}\right.  \tag{7.a}\\
& \left\{\begin{array}{l}
\mathrm{A}_{3}=\sqrt{\mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{4}\right)^{2}+\left(\frac{\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}-\mathrm{v}_{\mathrm{C}_{\mathrm{P}}}\left(\mathrm{t}_{4}\right)}{\mathrm{Z}_{3}}\right)^{2}} \\
\operatorname{tg}\left(\varphi_{3}\right)=\frac{\mathrm{Z}_{3} \cdot \mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{4}\right)}{-\left(\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}-\mathrm{v}_{\mathrm{C}_{\mathrm{P}}}\left(\mathrm{t}_{4}\right)\right)}
\end{array}\right.  \tag{7.b}\\
& \mathrm{v}_{\mathrm{C}_{\mathrm{P}}}\left(\mathrm{t}_{4}\right)=\mathrm{v}_{\mathrm{C}_{\mathrm{R} 0}}+\mathrm{V}_{\mathrm{i}}
\end{aligned} \omega_{3}=\frac{1}{\sqrt{\mathrm{C}_{\mathrm{P}} \cdot \mathrm{~L}_{\mathrm{R}}}}, ~ \begin{aligned}
& \mathrm{Z}_{3}=\sqrt{\frac{\mathrm{L}_{\mathrm{R}}}{\mathrm{C}_{\mathrm{P}}}}
\end{align*}
$$

At the same time the magnetizing current continues to decrease linearly following the relation (5) given in the previous interval.
6) Interval $t_{5}{ }^{-t} t_{6}$ (Fig.3f): After conduction of diode $\mathrm{D}_{1}$, the current in $\mathrm{L}_{\mathrm{R}}$ increases linearly with a slope proportional to $\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}$ :

$$
\begin{equation*}
\mathrm{i}_{\mathrm{R}}(\mathrm{t})=\mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{5}\right)+\frac{\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}}{\mathrm{~L}_{\mathrm{R}}} \cdot \mathrm{t} \tag{8}
\end{equation*}
$$

Note that $S_{1}$ must be turned on when $i_{R}$ is still negative in order to obtain a zero voltage turn on. At $\mathrm{t}_{6}$ current $\mathrm{i}_{\mathrm{R}}$ becomes equal to the magnetizing current $\mathrm{i}_{\mathrm{M}}$ (which was still linearly decreasing), and the freewheeling diode D stops conduction, initiating another switching cycle.

## B. Discontinuous Conduction Mode (DCM)

Fig. 4 reports the main waveforms of the same converter taken at $5 \%$ of rated power. The only difference in the operation sequence is that interval $\mathrm{t}_{3}-\mathrm{t}_{4}$ is terminated by the turn off of the freewheeling diode D while $\mathrm{S}_{2}$ is still on. After that, $\mathrm{L}_{\mathrm{R}}$ and $\mathrm{L}_{\mathrm{M}}$ resonate with $C_{R}$ until $S_{2}$ is turned off (instant $t_{5}$ ). Then, $L_{R}$ and $L_{M}$ together resonate with $\mathrm{C}_{\mathrm{P}}$, bringing its voltage to zero (instant $\mathrm{t}_{6}$ ) and initiating another switching cycle.


Fig. 4 - Main converter waveforms in DCM operation

## III. Extension to Isolated CUK and SEPIC TOPOLOGIES

The analysis done for the flyback converter can be easily extended also to Sepic and Cuk topologies.

## A. Sepic Converter

Consider the Sepic converter with active clamp shown in Fig.5a. As for the flyback converter, $\mathrm{L}_{\mathrm{R}}$ represents the transformer leakage inductance and $C_{P}$ the main switch output capacitance.

Assuming that energy transfer capacitor $\mathrm{C}_{1}$ has a negligible voltage ripple, it can be substituted by a voltage generator $\mathrm{V}_{\mathrm{C} 1}$ having the same value of the input voltage (which is the average voltage across $\mathrm{C}_{1}$ in steady state). Moving it across node A and rearranging we obtain the scheme drawn in Fig.5b. If the approximation shown in Fig. $5 c$ holds i.e. $\mathrm{L}_{1}, \mathrm{~L}_{\mathrm{M}} \gg \mathrm{L}_{\mathrm{R}}$, the scheme of Fig. 5 b coincides with that of Fig.1.


Fig. 5 - a) Sepic converter with active clamp; b) the same converter considering $\mathrm{C}_{1}$ as a voltage generator;
c) approximation which holds if $\mathrm{L}_{1}, \mathrm{~L}_{\mathrm{M}} \gg \mathrm{L}_{\mathrm{R}}$

## A. Cuk Converter

Let's consider now the Cuk converter with active clamp shown in Fig.6a in which the output filter capacitor and the load are substituted by a voltage generator $\mathrm{V}_{\mathrm{O}}$. Assuming that energy transfer capacitors $C_{1}$ and $C_{2}$ have a negligible voltage ripple, they can be substituted by voltage generators $\mathrm{V}_{\mathrm{C}}$ and $\mathrm{V}_{\mathrm{C}}$ having the same value of the input and output voltages respectively (which are the average voltages across $\mathrm{C}_{1}$ and $\mathrm{C}_{2}$ in steady state). Moving $\mathrm{V}_{\mathrm{C} 1}$ across node A and $\mathrm{V}_{\mathrm{C} 2}$ across node B and rearranging we obtain the scheme drawn in Fig.6b. If the approximation shown in Fig.6c holds i.e. $\mathrm{L}_{1}, \mathrm{~L}_{\mathrm{M}}, \mathrm{L}_{2} / \mathrm{n}^{2} \gg \mathrm{~L}_{\mathrm{R}}$, the scheme of Fig. 6b coincides with that of Fig.1.

It is worthy to note that also for the Cuk converter, the freewheeling diode average current is equal to the load current because the average current in $\mathrm{C}_{2}$ is zero in steady state.

## IV. Simplified Converter Analysis

As we can see from the waveforms of Figs. 2 and 4, the behavior of the flyback converter with active clamp is quite different respect to its hard-switched counterpart. However, using suitable approximations, it is possible to derive useful relations which can greatly simplify the design procedure.

b)

c)

Fig. 6 - a) Cuk converter with active clamp; b) the same converter considering $\mathrm{C}_{1}$ and $\mathrm{C}_{2}$ as voltage generators;
c) approximation which holds if $\mathrm{L}_{1}, \mathrm{~L}_{\mathrm{M}}, \mathrm{L}_{2} / \mathrm{n}^{2} \gg \mathrm{~L}_{\mathrm{R}}$

First of all, as far as the conversion ratio is concerned, we can observe that the voltage across the magnetizing inductance $\mathrm{L}_{\mathrm{M}}$ is equal to $\mathrm{V}_{\mathrm{i}} /(1+\beta)$ during the switch on-time, while during the switch off-time it changes from $\mathrm{V}_{\mathrm{op}}$ (interval $\mathrm{t}_{3}-\mathrm{t}_{4}$ ) to ${ }^{\mathrm{v}_{\mathrm{C}_{\mathrm{R}}}} /(1+\beta)$ during interval $\mathrm{t}_{2}-\mathrm{t}_{3}$ (and also $\mathrm{t}_{4}-\mathrm{t}_{5}$ in DCM ). Since ${ }^{v_{C}} \mathrm{C}_{R}$ value in these intervals is close to $\mathrm{V}_{\mathrm{op}}$ and considering that the duration of this interval is small compared to that of interval $t_{3}-t_{4}$, we have from the voltage balance on $\mathrm{L}_{\mathrm{M}}$ :

$$
\begin{align*}
& \frac{\mathrm{V}_{\mathrm{i}}}{1+\beta} \cdot \delta \mathrm{T}_{\mathrm{S}} \cong \mathrm{~V}_{\mathrm{op}} \cdot(1-\delta) \mathrm{T}_{\mathrm{S}} \\
& \mathrm{M}=\frac{\mathrm{V}_{\mathrm{op}}}{\mathrm{~V}_{\mathrm{i}}}=\frac{\delta}{1-\delta} \cdot \frac{1}{(1+\beta)} \tag{9}
\end{align*}
$$

which is almost the same of the usual flyback converter. Note that the same expression holds for both CCM and DCM operating modes.

The average current $\mathrm{I}_{\mathrm{M}}$ in the magnetizing inductance, assuming unity efficiency, is given by:

$$
\begin{equation*}
\mathrm{I}_{\mathrm{M}}=\mathrm{I}_{\mathrm{i}}+\mathrm{nI}_{\mathrm{o}}=\mathrm{nI}_{\mathrm{o}}(1+\mathrm{M}) \tag{10}
\end{equation*}
$$

while its current ripple amplitude is:

$$
\begin{equation*}
\Delta i_{M}=\frac{V_{i}}{L_{M} f_{S}} \cdot \frac{\delta}{(1+\beta)}=\frac{V_{o p}}{L_{M} f_{S}} \cdot \frac{1}{1+M(1+\beta)} \tag{11}
\end{equation*}
$$

where $\mathrm{f}_{\mathrm{S}}=1 / \mathrm{T}_{\mathrm{S}}$ is the switching frequency.
From (10) and (11) the peak current in the magnetizing inductance results:

$$
\begin{equation*}
\hat{\mathrm{i}}_{\mathrm{M}}=\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right)=\mathrm{I}_{\mathrm{M}}+\frac{\Delta \mathrm{i}_{\mathrm{M}}}{2}=\mathrm{nI} \cdot\left[1+\mathrm{M}+\frac{1}{\mathrm{k}[1+\mathrm{M}(1+\beta)]}\right] \tag{12}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathrm{k}=\frac{2 \mathrm{~L}_{\mathrm{M}} \mathrm{f}_{\mathrm{S}}}{\mathrm{R}_{\mathrm{op}}}=\frac{2 \mathrm{~L}_{\mathrm{M}} \mathrm{f}_{\mathrm{S}} \mathrm{n}^{2}}{\mathrm{R}_{\mathrm{o}}} \tag{13}
\end{equation*}
$$

is the usual adimensional parameter used in the flyback analysis.
In order to derive the value of load current at which the operating mode changes between CCM and DCM we must introduce some approximations. In particular, if in interval $\mathrm{t}_{3}-\mathrm{t}_{4}$ we consider $\beta \approx 0$ (which means $\mathrm{L}_{\mathrm{R}} \ll \mathrm{L}_{\mathrm{M}}$ ) and if we consider $\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{3}\right) \approx \mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right)$, then from (6) current $i_{R}$ and voltage ${ }^{v_{C}}{ }_{R}$ simplify to:

$$
\left\{\begin{array}{l}
\mathrm{i}_{\mathrm{R}}(\mathrm{t})=\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right) \cdot \cos \left(\omega_{2} \mathrm{t}\right)  \tag{14}\\
\mathrm{v}_{\mathrm{C}_{\mathrm{R}}}(\mathrm{t})=\mathrm{V}_{\mathrm{op}}+\mathrm{Z}_{2} \mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right) \cdot \sin \left(\omega_{2} \mathrm{t}\right)
\end{array}\right.
$$

Thus,

$$
\begin{equation*}
\mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{4}\right)=\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right) \cos \left(\omega_{2} \mathrm{~T}_{34}\right)=\mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right) \cos (\varepsilon) \tag{15.a}
\end{equation*}
$$

where

$$
\begin{equation*}
\varepsilon=\omega_{2} \mathrm{~T}_{34} \text { and } \mathrm{T}_{34}=\mathrm{t}_{3}-\mathrm{t}_{4} \approx(1-\delta) \mathrm{T}_{\mathrm{S}} . \tag{15.b}
\end{equation*}
$$

The discontinuous mode of operation is reached when at instant $t_{4}$ resonant inductor current $i_{R}$ and magnetizing current $i_{M}$ are equal. Thus:

$$
\begin{equation*}
\mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{4}\right)=\mathrm{I}_{\mathrm{M}}-\frac{\Delta \mathrm{i}_{\mathrm{M}}}{2} \tag{16}
\end{equation*}
$$

which gives the value of parameter Kcrit at the boundary between CCM and DCM:

$$
\begin{equation*}
\mathrm{k}_{\mathrm{crit}}=\frac{1-|\cos (\varepsilon)|}{1+|\cos (\varepsilon)|} \cdot \frac{1}{(1+\mathrm{M})^{2}} \tag{17}
\end{equation*}
$$

This means that the converter enters in DCM at a fraction $\alpha$ of the nominal output power, given by:

$$
\begin{equation*}
\alpha=\mathrm{R}_{\mathrm{op}} \cdot \frac{\mathrm{k}_{\text {crit }}}{2 \mathrm{~L}_{\mathrm{M}} \mathrm{f}_{\mathrm{S}}} \tag{18}
\end{equation*}
$$

## V. Zero Voltage Switching Conditions

The analysis reported in section III shows that in CCM operation only the energy stored in the leakage inductance at instant $t_{4}$ plays a role in discharging the mosfet parasitic capacitance $\mathrm{C}_{\mathrm{P}}$, while in DCM operation also the magnetizing inductance $\mathrm{L}_{\mathrm{M}}$ is involved in this process (instant $\mathrm{t}_{5}$ ). Thus, the condition to obtain zero voltage commutation depends on the operating mode.

## A. ZVS Condition in CCM

From the previous analysis, looking at the interval $\mathrm{t}_{4}-\mathrm{t}_{5}$, we can see that the voltage ${ }^{{ }^{2}} \mathrm{C}_{\mathrm{P}}$ reaches zero if the following inequality holds:

$$
\begin{equation*}
A_{3} Z_{3} \geq V_{i}+V_{o p} \tag{19}
\end{equation*}
$$

Using the same approximation for which (14) was derived, and assuming ${ }^{\mathrm{V}_{\mathrm{R} 0}} \approx \mathrm{~V}_{\mathrm{op}}$, (19) can be written as:

$$
\begin{equation*}
\mathrm{Z}_{3} \mathrm{i}_{\mathrm{M}}\left(\mathrm{t}_{1}\right)|\cos (\varepsilon)| \geq \mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}} \tag{20}
\end{equation*}
$$

which, rearranged, gives the following constraint on $\mathrm{L}_{\mathrm{R}}$ :

$$
\begin{equation*}
\sqrt{\frac{\mathrm{L}_{\mathrm{R}}}{\mathrm{C}_{\mathrm{P}}}} \geq \frac{2 \mathrm{~L}_{\mathrm{M}} \mathrm{f}_{\mathrm{S}}}{\mathrm{M}} \cdot \frac{(1+\mathrm{M})^{2}}{1+\mathrm{k}(1+\mathrm{M})^{2}} \cdot \frac{1}{|\cos (\varepsilon)|} \tag{21}
\end{equation*}
$$

The minimum value of $L_{R}$ is obtained for $\varepsilon=\pi$.
Note that if inductance $\mathrm{L}_{\mathrm{R}}$ is designed so as to meet the above inequality for $\mathrm{k}=0$, then the ZVS condition should be maintained in the entire load range.

## B. ZVS Condition in DCM

In order to derive the ZVS condition for this operating mode, the interval $\mathrm{t}_{5}-\mathrm{t}_{6}$ of Fig. 4 must be analyzed. The subcircuit corresponding to this interval is equal to that of Fig.3b. From this latter, the behavior of inductor currents and switch voltage is:

$$
\begin{align*}
& \left\{\begin{array}{l}
\mathrm{i}_{\mathrm{R}}(\mathrm{t})=\mathrm{i}_{\mathrm{M}}(\mathrm{t})=\mathrm{A}_{4} \cdot \sin \left(\omega_{4} \mathrm{t}+\varphi_{4}\right) \\
\mathrm{v}_{\mathrm{C}_{\mathrm{P}}}(\mathrm{t})=\mathrm{V}_{\mathrm{i}}-A_{4} \mathrm{Z}_{4} \cdot \cos \left(\omega_{4} \mathrm{t}+\varphi_{4}\right)
\end{array}\right.  \tag{22.a}\\
& \left\{\begin{array}{l}
\mathrm{A}_{4}=\sqrt{\mathrm{i}_{\mathrm{R}}\left(\mathrm{t}_{5}\right)^{2}+\left(\frac{\mathrm{v}_{\mathrm{C}_{\mathrm{R} 0}}}{\mathrm{Z}_{4}}\right)^{2}} \\
\operatorname{tg}\left(\varphi_{4}\right)=\frac{\mathrm{Z}_{4} \cdot i_{R}\left(\mathrm{t}_{5}\right)}{-\mathrm{v}_{\mathrm{C}_{\mathrm{R} 0}}}
\end{array}\right.  \tag{22.b}\\
& \omega_{4}=\frac{1}{\sqrt{\mathrm{C}_{\mathrm{P}} \cdot \mathrm{~L}_{\mathrm{M}}(1+\beta)}}  \tag{22.c}\\
& Z_{4}=\sqrt{\frac{L_{M}(1+\beta)}{C_{P}}} \tag{22.d}
\end{align*}
$$

In order for ${ }^{{ }^{2}} \mathrm{C}_{\mathrm{P}}$ to reach zero the following condition must be satisfied:

$$
\begin{equation*}
\mathrm{A}_{4} \mathrm{Z}_{4} \geq \mathrm{V}_{\mathrm{i}} \Rightarrow \mathrm{Z}_{4}^{2} \geq \frac{\mathrm{V}_{\mathrm{i}}^{2}-\mathrm{v}_{\mathrm{C}_{\mathrm{R} 0}}^{2}}{\mathrm{i}_{\mathrm{R}}^{2}\left(\mathrm{t}_{5}\right)} \tag{23}
\end{equation*}
$$

Considering that $i_{R}\left(\mathrm{t}_{5}\right)=\mathrm{I}_{\mathrm{M}^{-}-\Delta \mathrm{I}_{M}}$ 2, considering $\beta=0$ and using the approximation ${ }^{{ }^{\mathrm{C}}} \mathrm{R}_{\mathrm{R} 0} \approx \mathrm{~V}_{\mathrm{op}}$, the above constraint yields;

$$
\begin{equation*}
\sqrt{\frac{\mathrm{L}_{\mathrm{M}}}{\mathrm{C}_{\mathrm{P}}}} \geq \frac{2 \mathrm{~L}_{\mathrm{M}} \mathrm{f}_{\mathrm{S}}}{\mathrm{M}} \cdot \frac{(1+\mathrm{M}) \cdot \sqrt{1-\mathrm{M}^{2}}}{\left|\mathrm{k}(1+\mathrm{M})^{2}-1\right|} \tag{24}
\end{equation*}
$$

## VI. Design Guidelines

From the above analysis we are now able to derive a suitable design procedure which allows zero voltage commutation in all operating conditions.

## A. Magnetizing Inductance

The choice of the magnetizing inductance, besides the ZVS condition, influences both current and voltage stresses on the main switch. In fact, the switch current stress is equal to the peak inductor current reported in (12) while the switch voltage stress also depends on the peak magnetizing current due to the resonance between $C_{R}$
and $\mathrm{L}_{\mathrm{R}}$ during interval $\mathrm{t}_{3}-\mathrm{t}_{4}$. In fact, from the second of (14) it results:

$$
\begin{equation*}
\hat{\mathrm{v}}_{\mathrm{S}_{1}}=\mathrm{V}_{\mathrm{i}}+\mathrm{v}_{\mathrm{C}_{\mathrm{R}, \mathrm{MAX}}} \cong \mathrm{~V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}+\mathrm{Z}_{2} \cdot \hat{\mathrm{i}}_{\mathrm{M}} \tag{25}
\end{equation*}
$$

Both $Z_{2}$ and $\hat{\mathrm{i}}_{\mathrm{M}}$ depend on the magnetizing inductance value $\left(\mathrm{Z}_{2}\right.$ indirectly through (21)), but their product decreases for lower inductance values.

Once the desired current ripple $r_{i}$ on $L_{M}$ is chosen, the inductance is calculated by:

$$
\begin{equation*}
\mathrm{L}_{\mathrm{M}} \cong \frac{\mathrm{R}_{\mathrm{op}}}{2 \mathrm{f}_{\mathrm{S}}(1+\mathrm{M})^{2}} \cdot \frac{1}{\mathrm{r}_{\mathrm{i}}} \tag{26}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathrm{r}_{\mathrm{i}}=\frac{\Delta \mathrm{i}_{\mathrm{M}}}{2 \mathrm{I}_{\mathrm{M}}} \tag{27}
\end{equation*}
$$

We must remember that a low current ripple implies a high resonant inductor value in order to maintain the soft switching condition. Note that (26) is the same as eq. (6) in ref. [2].

## B. Resonant Inductance $L_{R}$

The value of inductance $\mathrm{L}_{\mathrm{R}}$ is chosen so as to guarantee the zero voltage commutation in all operating condition. Thus, the following procedure should be adopted:

- a suitable value of angle $\varepsilon$ defined by (15.b) is chosen;
- the value of critical parameter $\mathrm{k}_{\text {crit }}$ is found from (17);
- $\quad L_{R}$ value is calculated from (21) for $k=k_{c r i t}$ which represents the worst condition;
- the ZVS condition in DCM must be verified from (24) for $\mathrm{k}=\mathrm{k}_{\text {crit }}$.
As far as the choice of angle $\varepsilon$ is concerned, we can see from (21) that a value different from $\pi$ causes an increases of resonant inductor value needed to maintain the soft switching condition because it reduces the current, and thus the energy, available at instant $t_{4}$ (see Fig. 2). Thus it is convenient to set $\varepsilon=\pi$ in the worst condition, i.e. at maximum input voltage for which the magnetizing current is minimum.

As suggested in [2], the value of $\mathrm{L}_{\mathrm{R}}$ can be chosen also to limit the rate of change of freewheeling diode current during interval $\mathrm{t}_{5}$ $\mathrm{t}_{6}$ (see Fig.2).

## C. Resonant Capacitor $C_{R}$

The value of resonant capacitor $\mathrm{C}_{\mathrm{R}}$ is determined by the resonant period in the interval $t_{3}-t_{4}$, thus from the value of angle $\varepsilon$ chosen in the previous subsection. From (6.c), (15.b) and (9) and using the approximations $\beta=0$ and $\mathrm{C}_{\mathrm{P}} \ll \mathrm{C}_{\mathrm{R}}$ we can write:

$$
\begin{equation*}
\mathrm{C}_{\mathrm{R}} \approx \frac{1}{\mathrm{~L}_{\mathrm{R}}\left[\varepsilon \mathrm{f}_{\mathrm{S}}(1+\mathrm{M})\right]^{2}} \tag{28}
\end{equation*}
$$

D. Main Switch $S_{1}$

Switch $S_{1}$ is chosen on the basis of its current and voltage stresses. These latter can be found by using (12) and (25) respectively.
E. Auxiliary Switch $S_{2}$

Also switch $S_{2}$ is chosen on the basis of its maximum ratings. As far as its voltage stress is concerned, it is easily verified that it is given by the following equation:

$$
\begin{equation*}
\hat{\mathrm{v}}_{\mathrm{S}_{2}}=\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{C}_{\mathrm{R} 0}} \tag{29}
\end{equation*}
$$

while its current stress is the same of the main switch because the current flowing in it coincides practically with resonant current $\mathrm{i}_{\mathrm{R}}$ during interval $t_{2}{ }^{-t_{4}}$.

## F. Switch Drive Signals

In order to obtain the correct converter behavior, suitable switch drive signal must be generated; in particular, dead times between turn off of $S_{1}$ and turn on of $S_{2}$, and vice versa, are needed in order to provide a complete charge and discharge of parasitic capacitance $C_{P}$. From (7), the dead time $t_{d 1}$ between turn off of $S_{2}$ and turn on of $S_{1}$ in CCM can be approximated by

$$
\begin{equation*}
\mathrm{t}_{\mathrm{d} 1} \approx \frac{\pi}{2 \omega_{3}}=\frac{\pi}{2} \cdot \sqrt{\mathrm{~L}_{\mathrm{R}} \mathrm{C}_{\mathrm{P}}} \tag{30}
\end{equation*}
$$

while in $D C M L_{R}$ in (30) should be replaced by $L_{R}+L_{M}$. In practice, in DCM the discharge of $C_{P}$ occurs in an interval of time much shorter than one half of resonant period, thus the suitable delay time lays between the two values.

The dead time $t_{d 2}$ between turn off of $S_{1}$ and turn on of $S_{2}$ can be derived considering interval $t_{1}{ }^{-} t_{2}$ and (2) in the worst case which occurs when the peak magnetizing inductor current is minimum, i.e.

$$
\begin{gather*}
\hat{\mathrm{i}}_{\mathrm{M}}=\frac{\Delta \mathrm{i}_{\mathrm{M}}}{2}, \text { with } \Delta \mathrm{i}_{\mathrm{M}} \text { given by }(11) \\
\mathrm{t}_{\mathrm{d} 2}=\frac{2 \mathrm{C}_{\mathrm{P}}\left(\mathrm{~V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}\right)}{\Delta \mathrm{i}_{\mathrm{M}}} \tag{31}
\end{gather*}
$$

## G. Capacitor $C_{P}$

This capacitance usually represents the sum of the parasitic output capacitances of $S_{1}$ and $S_{2}$ and of the freewheeling diode (reported to the primary side of the transformer). In some applications, it could be convenient to increase this capacitance by adding an external capacitor in order to limit the maximum voltage rate of change across the active devices. This provision reduces both conducted and radiated disturbances and can be necessary when IGBT are used instead of mosfets in order to reduce their turn off losses due the tail current phenomenon.

When an external capacitor is added, it is interesting to see the variation of the normalized switch voltage stress $\hat{\mathrm{v}}_{\mathrm{S}_{\mathrm{N}}}=\hat{\mathrm{v}}_{\mathrm{S}_{1}} /\left(\mathrm{V}_{\mathrm{i}}+\mathrm{V}_{\mathrm{op}}\right)$ as a function of the maximum $\mathrm{dv} / \mathrm{dt}$ allowed as shown in Fig.7. which reports this behavior for three different values of magnetizing current ripple $r_{i}$.

As we can see, the voltage stress increases rapidly as the maximum dv/dt is reduced. This is due to the increase of characteristic impedance $Z_{2}$ caused by the contemporary increase of $\mathrm{L}_{\mathrm{R}}$, needed to maintain the soft switching condition, and decrease of $\mathrm{C}_{\mathrm{R}}$.


Fig. 7 - Normalized switch voltage stress versus maximum dv/dt for three different values of magnetizing current ripple

It is worthy to note that the behavior reported in Fig. 7 is only indicative of the phenomenon because, as $\mathrm{C}_{\mathrm{P}}$ increases, the dead times $\mathrm{t}_{\mathrm{d} 1}$ and $\mathrm{t}_{\mathrm{d} 2}$ increase becoming non negligible respect to the switching period and thus affecting the validity of the above relations.

## VII. EXPERIMENTAL RESULTS

A prototype, based on flyback topology, was built and tested.
Specifications and parameter values are listed in Table I. The main components used are:
$S_{1}, S_{2}$ : IRF 640R, D: RUR 1520
$\mathrm{L}_{\mathrm{M}}$ : Core E3213A, $\mathrm{N}_{1}=\mathrm{N}_{2}=12$ turns litz wire
$\mathrm{L}_{\mathrm{R}}$ : Core $\mathrm{P} 22 / 13, \mathrm{~N}=3$ turns litz wire
The control strategy is the peak current control implemented with the UC3823A integrated controller. The two drive signals are generated by a dual power mosfet driver (TC428) commanded by the output of the UC3823A through two delay networks.

Figs.8a) and b) show resonant inductor current $\mathrm{i}_{\mathrm{LR}}$, main switch drain-to-source voltage ${ }^{\text {v DS1 }}$ and gate-to-source voltage ${ }^{\mathrm{v}} \mathrm{GS} 1$ waveforms at $100 \%$ and about $10 \%$ of rated power respectively. As we can see, soft-switching condition is maintained also at the boundary between discontinuous and continuous conduction mode, which is the most critical point. In fact, at lower output currents, the soft-switching condition is ensured by the energy stored also in the transformer magnetizing inductance. Observe that the switch voltage waveform is very smooth with just a small ringing at the beginning of the off interval due to layout-induced parasitic inductance. Thus, the electromagnetic noise generated is reduced.

The voltage and current stresses in the main switch are 172 V and 5.4 A respectively. The behavior of the overall efficiency including control stage as a function of the load current is reported in Fig.9. As we can see, it remains above $83 \%$ for output currents greater than $50 \%$ of the nominal value.

Table I - Converter specifications and parameters

| $\mathrm{V}_{\mathrm{i}}=100 \mathrm{~V}$ | $\mathrm{~V}_{\mathrm{O}}=48 \mathrm{~V}$ | $\mathrm{f}_{\mathrm{S}}=300 \mathrm{kHz}$ | $\mathrm{R}_{\mathrm{O}}=24 \Omega$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{r}_{\mathrm{i}}=0.75$ | $\varepsilon=7 \pi / 6$ | $\mathrm{~L}_{\mathrm{M}}=25 \mu \mathrm{H}$ | $\mathrm{L}_{\mathrm{R}}=3.8 \mu \mathrm{H}$ |
| $\mathrm{C}_{\mathrm{O}}=100 \mu \mathrm{~F}$ | $\mathrm{C}_{\mathrm{R}}=100 \mathrm{nF}$ | $\mathrm{C}_{\mathrm{P}}=0.6 \mathrm{nF}$ | $\mathrm{n}=1$ |



Fig. 8 - Resonant inductor current ${ }^{\mathrm{L}}$ LR, main switch drain-to-source voltage $\mathrm{V}_{\mathrm{DS} 1}$ [50V/div] and main switch gate-to-source voltage $\mathrm{V}_{\mathrm{GS} 1}$ [5V/div]; a) $\mathrm{i}_{\mathrm{LR}}[5 \mathrm{~A} / \mathrm{div}], \mathrm{I}_{\mathrm{O}}=2 \mathrm{~A} ; \mathrm{b}$ ) ${ }^{\mathrm{i}}{ }_{\mathrm{LR}}[2 \mathrm{~A} / \mathrm{div}], \mathrm{I}_{\mathrm{O}}=0.18 \mathrm{~A}$


Fig. 9 - Overall converter efficiency vs. load current CONCLUSIONS

A flyback converter with active voltage clamp was analyzed reviewing its behavior for both CCM and DCM operation modes.

Suitable design criteria are given which allow soft commutation of all devices from no load to full load. Moreover, It is shown that
the analysis done is also applicable to other insulated DC/DC topologies like Cuk and Sepic.

The power stage is very simple and reliable and the control does not need particular provisions respect to standard PWM controller.

Experimental results of a $100 \mathrm{~W}-300 \mathrm{kHz}$ converter confirmed the theoretical forecasts showing good performances.

## References

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